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Microwave Dielectric Properties of Temperature-Stable Zircon-Type (Bi, Ce)VO₄ Solid-Solution Ceramics

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Abstract

In the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0 \leq x \leq 1$) system, we found that the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0 \leq x \leq 0.1$) belongs to the monoclinic scheelite phase and the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq x \leq 1$) belongs to the tetragonal zircon phase, while the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.1 < x < 0.7$) belongs to the mixed phases of both monoclinic scheelite and tetragonal zircon structure. Interestingly, two components with near-zero temperature coefficient of resonant frequency (TCF) appeared in this system. In our previous work, a near-zero TCF of $\sim +15 \text{ ppm}/^\circ\text{C}$ was obtained in a $(\text{Bi}_{0.75}\text{Ce}_{0.25})\text{VO}_4$ ceramic with a permittivity (ϵ_r) of ~ 47.9 , a Qf ($Q = \text{quality factor} = 1/\text{dielectric loss}$; $f = \text{resonant frequency}$) value of $\sim 18,000 \text{ GHz}$ (at 7.6 GHz). Furthermore, in the present work, another temperature stable microwave dielectric ceramic was obtained in $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ composition sintered at 950°C and exhibits good microwave dielectric properties with a ϵ_r of

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~ 11.9, a Qf of ~ 22,360 GHz (at 10.6 GHz), a near-zero TCF of ~ +6.6 ppm/°C. The results indicate that this system might be an interesting candidate for microwave device applications.

KEYWORDS: low temperature; X-ray methods; solid solutions; dielectric properties

1. INTRODUCTION

Microwave dielectric materials play a key role in global society, with a wide range of applications straddling terrestrial and satellite communications, including Internet of Things (IoT), software radio, GPS, DBS TV, environmental monitoring via satellite, etc. The temperature coefficient of resonance frequency (TCF) is one of the important parameters of microwave dielectric ceramics. The dielectric resonator generally uses the frequency of a certain vibration mode of dielectric ceramic as its center frequency. Therefore, in order to eliminate the temperature drift of the resonant frequency characteristic of the resonator, it is necessary to make the TCF close to zero.¹⁻⁷

In 2000, Valant and Suvorov⁸ first reported that BiVO₄ ceramic not only sinter at a low temperature (< 900 °C) but also exhibit excellent microwave dielectric properties, with a permittivity (ϵ_r) of ~ 68, a Qf (Q = quality factor = 1/dielectric loss; f = resonant frequency) value of ~ 6,500 GHz, a negative TCF of ~ -260 ppm/°C. However, the large negative temperature coefficient of resonant frequency will limit its application in microwave devices, thus, further studies to adjust the TCF of BiVO₄ ceramic to near zero is of great interest. In

our previous work,⁹ the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($x \leq 0.6$) ceramics were prepared via a solid-state reaction method and a near-zero temperature coefficient of resonance frequency (+15 ppm/°C) was obtained in a $(\text{Bi}_{0.75}\text{Ce}_{0.25})\text{VO}_4$ ceramic with a permittivity of ~ 47.9 , a Qf value of $\sim 18,000$ GHz. However, we have noticed that the $(\text{Bi}_{0.4}\text{Ce}_{0.6})\text{VO}_4$ has a TCF = +173 ppm/°C and the microwave dielectric properties of zircon-structured CeVO_4 were reported by Zuo et al.¹⁰ with a permittivity of ~ 12.3 , a Qf value of $\sim 41,460$ GHz, and a negative TCF of ~ -34.4 ppm/°C. Furthermore, from the pseudo phase diagram of the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ system reported in the literature,^{9, 11} $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ should be a tetragonal zircon phase between $x = 0.6 \sim 1$. Therefore, there is a possibilities to design a temperature stable microwave dielectric ceramic in $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.6 < x < 1$) solid solutions.

In this investigation, the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq x \leq 0.95$) solid solution ceramics were synthesized by a conventional solid phase reaction method to prepare a temperature-stable microwave dielectric ceramic with near- zero TCF. The structural evolution, sintering behaviors, microstructures, microwave dielectric properties were investigated in detail.

2. EXPERIMENTAL PROCEDURES

Proportionate amounts of reagent-grade starting materials of CeO_2 (99%), Bi_2O_3 (99%), V_2O_5 (99%) were measured based on the stoichiometric formulation $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq x \leq 0.95$). Details of the processing can be found in our previous work.⁹ Samples were sintered in the temperature range $870 \sim 970$ °C for 2 h. The samples were investigated using room-temperature X-ray diffraction (XRD) with Cu K α radiation (Rigaku D/MAX-2400 X-ray diffractometry, Tokyo, Japan). The microstructure images of the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq$

$x \leq 0.95$) ceramics were obtained by scanning electron microscopy (FEI Quanta F250). The SAED and HRTEM images were measured using a Transmission electron microscope (JEM-2100 Electron Microscope, Japan) operated at 200 kV. The surface properties of the $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic was characterized using X-ray photoelectron spectroscopy (XPS) on a Thermo Fisher ESCALAB Xi+. The infrared reflectivity spectra and Raman spectra were taken by a Bruker IFS 66v FTIR spectrometer (NSRL, Hefei, China) and a LabRAM HR Evolution with laser excitation at 532 nm, respectively. The permittivity and Qf value of $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq x \leq 0.95$) ceramics were obtained using a network analyser (8720ES, Agilent, Palo Alto, CA). Since the TCF value was tested at the temperature ranging from 25 °C to 85 °C, the instruments used for the test were network analyzer and temperature chamber (Delta 9023, Delta Design, Poway, CA). The formula for calculating TCF is as follows:

$$TCF(\tau_f) = \frac{f_T - f_{T_0}}{f_{T_0} \times (T - T_0)} \times 10^6 \quad (1)$$

where, f_{T_0} and f_T is the resonant frequencies at temperatures T_0 and T , respectively.

3. RESULTS AND DISCUSSION

Figure 1a shows the X-ray diffraction patterns of $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq x \leq 0.95$) ceramics sintered at their optimum temperatures. It can be seen that the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq x \leq 0.95$) were crystallized in a zircon-type tetragonal structure and no second phase is found, which indicates that the zircon-type tetragonal structure (space group $I4_1/amd$) is stable in the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq x \leq 0.95$) ceramics. In order to study the crystal structure of $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic in more detail, refinements were performed by using GSAS

software based on the X-ray diffraction data recorded on $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic powders.

The refinement results indicate that the $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic is a tetragonal zircon

structure with a space group $I4_1/amd$ (No. 141) and the lattice parameter are $a = b =$

$7.3777(4) \text{ \AA}$, $c = 6.4809(1) \text{ \AA}$. Figure 1b shows the calculated and measured X-ray diffraction

patterns, and the illustration is schematic crystal structure of the $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic. In

addition, the refined atomic fractional coordinates are listed in Table 1. The goodness of fit of

refinement ($R_{\text{wp}}/R_{\text{exp}}$) is 1.9. Furthermore, we also used transmission electron microscope

(TEM) to further analyze the structure of $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic. The selected area

electron diffraction (SAED) pattern and high-resolution transmission electron microscopy

(HRTEM) images of $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic recorded along the $[010]$ zone axis are

presented in Figure 1c, d. The HRTEM image demonstrates that the characteristic spacings of

the (200) and (101) lattice planes of $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic are 0.36 nm and 0.49 nm,

respectively, which correspond well with the refined lattice parameters. The schematic

diagram of crystal structure can be well matched with the HRTEM image, and the small dots

shown in Figure 1d should be oxygen atoms. All of these results are well matched, which

fully confirms that $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic belongs to the tetragonal zircon structure with a

space group $I4_1/amd$ (No. 141).

The surface micro-morphology of the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq x \leq 0.95$) ceramics were

studied by the scanning electron microscopy (SEM). The SEM images of thermally etched

surface of the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq x \leq 0.95$) ceramics sintered at optimal temperatures are

shown in Figure 2a ~ f. Dense and homogeneous microstructures with almost no pores could

be revealed in all compositions. Furthermore, the sintering temperatures of BiVO_4 and

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CeVO₄ ceramics are 820 °C and 950 °C, respectively.^{8, 10} So the optimum sintering temperature of (Bi_{1-x}Ce_x)VO₄ (0.7 ≤ x ≤ 0.95) ceramics increased from 890 °C to 950 °C as x increased. This results indicates that the sintering temperatures of CeVO₄ ceramic can be lowered by the formation of solid solution and similar phenomena also exist in many other solid solution systems, such as (1-x)BiVO₄-xLaNbO₄, (1-x)BiVO₄-xLa_{2/3}MoO₄, (Ca_{1-x}Bi_x)(Mo_{1-x}V_x)O₄ and (1-x)BiVO₄-x(Ag_{0.5}Bi_{0.5})MoO₄ etc.¹²⁻¹⁵

A temperature stable microwave dielectric ceramic with good microwave dielectric properties could be obtained at x = 0.95 (as discussed in Figure 5 below). Therefore, it is necessary to study the (Bi_{0.05}Ce_{0.95})VO₄ in more detail. Raman spectroscopy is a good means to analyse the vibration modes of structural units such as [VO₄]. Therefore, Raman spectroscopic investigation was also performed on (Bi_{0.05}Ce_{0.95})VO₄ ceramic. Pure CeVO₄ belongs to a tetragonal zircon structure with a space group I4₁/amd and a point group D_{4h} (4/mmm). It is predicted by group theory that it has twelve distinct Raman active modes:

$$\Gamma = 2A_{1g} + 4B_{1g} + B_{2g} + 5E_g \quad (2)$$

These vibration modes can be divided into two categories: internal models and external models. In CeVO₄, the vibration of the internal models is derived from the intermolecular vibration of the VO₄³⁻. The vibration of each VO₄³⁻ anion is associated with other VO₄³⁻ ions in the same unit cell, resulting in abundant multi-vibration structure. One of the peaks shown in the spectrum may be due to the mixing of the vibrational compositions of different types of molecules. In order to observe the lattice vibration modes, the standard Gaussian-Lorentzian model was used to fit the Raman spectrum of (Bi_{0.05}Ce_{0.95})VO₄ ceramic

to obtain the specific position information of all peaks. The fitted Raman spectra are presented in Figure 3a as red solid lines. As shown in Figure 3a, the strongest mode at 851 cm^{-1} is assigned as A_{1g} symmetric stretch (ν_1). The 777, 458, and 367 cm^{-1} Raman modes are assigned as E_g asymmetric stretch (ν_3), E_g+B_{2g} deformation (ν_4), and $A_{1g}+B_{1g}$ deformation (ν_2), respectively. The B_{2g} mode at 254 cm^{-1} can be interpreted as symmetric bending of the $[\text{VO}_4]$ tetrahedron. Furthermore, the external modes appear at frequencies below 250 cm^{-1} . All vibration modes are consistent with the results reported in the previous literature.^{9, 16, 17}

The four normal vibration modes of tetrahedral $[\text{VO}_4]$ are shown in Figure 3b. In addition, Figure 3c shows the calculated and measured infrared reflectivity spectra of $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic. The bands at 880, 788, and 259 cm^{-1} can be assigned as A_{2u} asymmetric stretch (ν_3), E_u asymmetric stretch (ν_3), and E_u deformation (ν_2), respectively. The weak band at 441 cm^{-1} is assigned as E_u+A_{2u} deformation (ν_4). While the band at 176 cm^{-1} can be assigned to the Bi–O band (E_u). The infrared spectra of $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic is similar to the reports in the literature.^{9, 17, 18} In order to investigate the intrinsic microwave dielectric properties in detail, the infrared reflectivity spectra of $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic were analyzed using a classical harmonic oscillator model:

$$\varepsilon^*(\omega) - \varepsilon(\infty) = \sum_{j=1}^n \frac{(z_j e)^2 / m_j V_j \varepsilon_0}{\omega_{Tj}^2 - \omega^2 - j\gamma_j \omega} \quad (3)$$

where, z_j is the equivalent charge of the j^{th} vibration mode, m_j is the equivalent mass of the j^{th} vibrational mode, V_j is the equivalent unit volume of the j^{th} vibrational mode, γ_j is the damping coefficient of the j^{th} vibrational mode, ω_{Tj} is the angular frequency of the transverse optical modes of the j^{th} mode of vibration, and n is the number of transverse

phonon modes. The relation between complex reflectivity $R(\omega)$ and dielectric constant can be written as follows:

$$R(\omega) = \left| \frac{1 - \sqrt{\varepsilon^*(\omega)}}{1 + \sqrt{\varepsilon^*(\omega)}} \right|^2 \quad (4)$$

Figure 3c shows the fitted infrared reflectivity values and the complex permittivities, and Table 2 lists the relevant parameters. It can be seen that the measured values (using the TE₀₁₈ method) are very close to the calculated dielectric constant and dielectric loss values, indicating that in the microwave region the dielectric polarization is mainly caused by phonon absorption in the infrared region.

X-ray photoelectron spectroscopy (XPS) was used to characterize the chemical composition and oxidation state of each element in the (Bi_{0.05}Ce_{0.95})VO₄ ceramic. The XPS spectrum of (Bi_{0.05}Ce_{0.95})VO₄ sample with the survey and high-resolution spectrum of all elements were shown in Figure 4a ~ e. It is seen that the survey spectrum (Figure 4a) revealed the presence of Bi, Ce, V, and O in (Bi_{0.05}Ce_{0.95})VO₄. In the Ce 3d spectrum (Figure 4c), The peaks at 881.8 and 885.9 eV correspond to the Ce 3d_{5/2} transitions, while the peaks at 900.2 and 904.3 eV belong to the Ce 3d_{3/2} binding energy. The binding energies at 885.9 and 904.3 eV represent the 3d¹⁰4f¹ initial electronic state corresponding to Ce³⁺ and no peak associated with Ce⁴⁺ was found, which is consistent with the results of CeVO₄ reported in the literature.^{19, 20} The V 2p spectrum exhibits two contributions, 2p_{3/2} and 2p_{1/2} (resulting from the spin-orbit splitting), located at respectively 517.1 and 524.6 eV, which can be matched well with V₂O₅.²¹ The Bi 4f_{7/2} peak at 159.0 eV agreed with the peak position of BiVO₄ in the literature,²² and the peak located at 530.0 eV can be assigned to O 1s of O²⁻.²³ These results

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indicated that the oxidation states of Bi, Ce and V in the $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ sample were +3, +3 and +5, respectively.

According to our previous work,⁹ we found that the TCF of BiVO_4 can be adjusted from -260 ppm/ $^{\circ}\text{C}$ to near zero (+15 ppm/ $^{\circ}\text{C}$) when the substitution amount of Ce is 0.25 in the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0 \leq x \leq 0.6$) ceramics. However, we also noticed that the TCF of $(\text{Bi}_{0.4}\text{Ce}_{0.6})\text{VO}_4$ ceramics is +173 ppm/ $^{\circ}\text{C}$, while Zuo et al. reported that the TCF of CeVO_4 ceramic is -34.4 ppm/ $^{\circ}\text{C}$.¹⁰ Therefore, another component with a near-zero TCF will appear in $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.6 < x < 1$) ceramics, which prompted us to find such a temperature-stable microwave dielectric ceramic. To better understand how microwave dielectric properties change with the substitution of Ce for Bi, the microwave dielectric properties ϵ_r , Qf and TCF of the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0 \leq x \leq 1$) ceramics sintered at their respective optimum temperatures as a function of x value are presented in Figure 5a ~ c. Among them, the data of $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0 \leq x \leq 0.6$) ceramics is taken from our previous work,⁹ and the data of $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq x \leq 0.95$) ceramics comes from this work. In addition, the data of CeVO_4 ceramic is taken from the literature reported by Zuo et al.¹⁰ In order to make clear the changes of ϵ_r and Qf with x value, we performed a linear fit on them separately. It can be seen that the ϵ_r of samples decreases linearly with the increase of x, mainly because the permittivity depends on the molecular polarizability. It has been reported in previous literature that the ionic polarizability of Ce^{3+} is smaller than that of Bi^{3+} in the structural environment of monoclinic scheelite and tetragonal zircon phase, thus resulting in the permittivity of $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0 \leq x \leq 1$) decreases with the cerium content.⁹ As plotted in Figure 5b, the Qf values decreases roughly with the increase of x. In other words, the Qf

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values are inversely proportional to the permittivities. It is due to the intrinsic dielectric loss caused by absorptions of phonon oscillation in the lattice is proportional to the ε_r and the optimum Qf value depends on the intrinsic dielectric loss based on the classical harmonic oscillator model. Therefore, the Qf value of $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0 \leq x \leq 1$) ceramics shows an increasing trend. At microwave region, the Qf value has inverse relation to permittivity ($Q \times f \approx \frac{(ze)^2/mV\varepsilon_0}{2\pi\gamma \times (\varepsilon'(\omega) - \varepsilon(\infty))}$).²⁴ This relation has successfully explained the relationship between the Qf and the permittivity of many scheelite solid solutions such as $(\text{Bi}_{1-x}\text{Y}_x)\text{VO}_4$, $[(\text{Li}_{0.5}\text{Bi}_{0.5})_{1-x}\text{Ca}_x]\text{MoO}_4$ and $(\text{Ca}_{1-x}\text{Bi}_x)(\text{Mo}_{1-x}\text{V}_x)\text{O}_4$.^{15, 18, 24} Moreover, the facts show that the same qualitative relationship is also suitable for $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0 \leq x \leq 1$) ceramics. As presented in Figure 5c, the TCF first increases from a negative value to a positive value, and then decreases from a positive value to a negative value. Therefore, two component points with near zero TCF appear in $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0 \leq x \leq 1$) ceramics, which are $x = 0.25$ and $x = 0.95$, respectively, indicating that the substitution of Ce for Bi in BiVO_4 is an effective method to adjust TCF to near zero. To understand the effect of sintering temperature on the TCF values of these two component points, the TCF of $(\text{Bi}_{0.75}\text{Ce}_{0.25})\text{VO}_4$ and $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramics as a function of sintering temperature were shown in Figure 6. It can be seen that the TCF of the $(\text{Bi}_{0.75}\text{Ce}_{0.25})\text{VO}_4$ and $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramics tends to decrease first and then remain stable as the sintering temperature increases. In order to better understand the crystal structure of the two components with near zero TCF, the pseudo phase diagram of $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0 \leq x \leq 1$) ceramics and the associated schematic crystal structures of each region are shown in Figure 5d, e. It can be seen that the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0 \leq x \leq 0.1$) belongs to the monoclinic scheelite phase and the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq x \leq 1$) belongs to the

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tetragonal zircon phase, while the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.1 < x < 0.7$) belongs to the mixed phases of both monoclinic scheelite and tetragonal zircon structure. Thus, the two components with near zero TCF belong to the mixed phases and the tetragonal zircon phase, respectively. It is important to note that the $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic sintered at 950 °C can exhibit good microwave dielectric properties, with a $\epsilon_r \sim 11.9$, a $Q_f \sim 22,360$ GHz (at 10.6 GHz) and a $\text{TCF} \sim +6.6$ ppm/°C. In the microwave region, polarizability is the sum of both ionic and electronic components. Shannon²⁵ suggested that molecular polarizability (α) of complex substances maybe estimated by summing α of the constituent ions which for $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ is:

$$\alpha_{\text{Bi}_{0.05}\text{Ce}_{0.95}\text{VO}_4} = 0.05\alpha_{\text{Bi}^{3+}} + 0.95\alpha_{\text{Ce}^{3+}} + \alpha_{\text{V}^{5+}} + 4\alpha_{\text{O}^{2-}} \approx 17.11 \text{ \AA}^3 \quad (5)$$

where the ionic polarizabilities of Bi^{3+} , Ce^{3+} , V^{5+} and O^{2-} are 6.12 \AA^3 , 6.15 \AA^3 , 2.92 \AA^3 and 2.01 \AA^3 , respectively.²⁵ Considering the Clausius–Mosotti relation,²⁶

$$\epsilon_{\text{meas}} = \frac{3V + 8\pi\alpha}{3V - 4\pi\alpha} \Rightarrow \alpha = \frac{3V(\epsilon_{\text{meas}} - 1)}{4\pi(\epsilon_{\text{meas}} + 2)} \approx 16.51 \text{ \AA}^3 \quad (6)$$

where V is the cell volume ($352.76/4 = 88.19 \text{ \AA}^3$), the molecular polarizability may be obtained from ϵ_r to give $\sim 16.51 \text{ \AA}^3$ which is similar to the calculated value (17.11 \AA^3) based on Shannon's additive rule. In addition, the TCF can be defined as follows:

$$\text{TCF} = -\left(\alpha_l + \frac{1}{2}\tau_\epsilon\right) \quad (7)$$

where α_l is the thermal expansion coefficient and τ_ϵ is the temperature coefficient of dielectric constant. According to the thermal expansion data (as shown in Figure 7), the α_l of $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic is +3.9 ppm/°C between 25 and 85 °C. Hence, it can be calculated

according to equation (7) that τ_ϵ of the $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic is $-21.0 \text{ ppm}/^\circ\text{C}$. In addition, using the Clausius–Mosotti relation, Bosman and Havinga²⁷ derived an expression for τ_ϵ at constant pressure as follows:

$$\tau_\epsilon = \frac{1}{\epsilon} \left(\frac{\partial \epsilon}{\partial T} \right)_p = \frac{(\epsilon - 1)(\epsilon + 2)}{\epsilon} (A + B + C) = \left(\epsilon - \frac{2}{\epsilon} + 1 \right) (A + B + C) \quad (8)$$

$$A = -\frac{1}{3V} \left(\frac{\partial V}{\partial T} \right)_p, B = \frac{V}{\alpha_m} \left(\frac{\partial \alpha_m}{\partial V} \right)_T \cdot \frac{1}{3V} \left(\frac{\partial V}{\partial T} \right)_p, C = \frac{1}{3\alpha_m} \left(\frac{\partial \alpha_m}{\partial T} \right)_V$$

The value of the A+B term is about $6 \text{ ppm}/^\circ\text{C}$. Moreover, term C generally lies between -1 and $\sim -10 \text{ ppm}/^\circ\text{C}$ and represents the direct dependence of the polarizability on temperature.

According to equation (8), the $(A + B + C)$ value of the $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic can be calculated to $-1.6 \text{ ppm}/^\circ\text{C}$, which is within an acceptable range.

4. CONCLUSIONS

In the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq x \leq 0.95$) ceramics, tetragonal zircon-type structured solid solution was formed. As x increased from 0.7 to 0.95, the sintering temperature of $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq x \leq 0.95$) ceramics increased from 890 to $950 \text{ }^\circ\text{C}$, while the TCF decreased linearly from $+176.3$ to $+6.6 \text{ ppm}/^\circ\text{C}$. Importantly, the $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic sintered at $950 \text{ }^\circ\text{C}$ can exhibit good microwave dielectric properties, with a $\epsilon_r \sim 11.9$, a $Qf \sim 22,360 \text{ GHz}$ (at 10.6 GHz) and a $\text{TCF} \sim +6.6 \text{ ppm}/^\circ\text{C}$. It is proved that a temperature-stable microwave dielectric ceramic with zircon-type tetragonal phase can be obtained by the substitution of Ce for Bi in BiVO_4 . This work presents a novel method to modify the TCF of

BiVO₄-type materials. (Bi_{0.05}Ce_{0.95})VO₄ ceramic have a good potential for microwave equipment and devices with near zero temperature coefficient of resonance frequency.

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List of Table and Figure Captions:

Table 1 Refined atomic fractional coordinates from XRD data of the $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic and the cell parameters are $a = b = 7.3777(4) \text{ \AA}$, $c = 6.4809(1) \text{ \AA}$ with a space group $I4_1/amd$ (No. 141).

Table 2 Phonon parameters obtained from the fitting of the infrared reflectivity spectra of $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic.

FIGURE 1 (a) X-ray diffraction patterns of the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq x \leq 0.95$) ceramics sintered at optimal temperatures. (b) The experimental (circle) and calculated (line) X-ray powder diffraction profiles for $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ sample sintered at 950°C (The short vertical lines below the patterns mark the positions of Bragg reflections. The bottom continuous line is the difference between the observed and calculated intensity.). Insets is the schematic crystal structure of $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$. (c, d) SAED pattern and HRTEM image of $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ viewed along the $[010]$ zone axis.

FIGURE 2 (a ~ f) SEM images of thermally etched surface of the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0.7 \leq x \leq 0.95$) ceramics sintered at optimal temperatures.

FIGURE 3 (a) The experimental (circle) and calculated (black solid line) Raman spectra of $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic sintered at 950 °C. (The short dot lines represent the Gaussian-Lorentzian fitting results). (b) Four normal vibration modes of tetrahedral $[\text{VO}_4]$. (c) Measured and calculated infrared reflectivity spectra (solid line for fitting values and circle for measured values) and fitted complex dielectric spectra of $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic (square and circles are experimental in microwave region).

FIGURE 4 XPS spectra of $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic with (a) the survey spectrum and high-resolution spectra of (b) Bi 4f, (c) Ce 3d, (d) V 2p and (e) O 1s. (all peaks have been calibrated with respect to the C 1s peak at 284.8 eV).

FIGURE 5 The microwave dielectric properties ϵ_r (a), Qf (b) and TCF (c) of $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ ($0 \leq x \leq 1$) ceramics sintered at optimal temperatures as a function of x value. (d) The pseudo phase diagram of the $(\text{Bi}_{1-x}\text{Ce}_x)\text{VO}_4$ system. The red dotted line represent the zircon phase contents in samples calculated by Fullprof software (Detailed information on the calculation method is described in the literature⁹). (e) The associated schematic crystal structures of each region in pseudo phase diagram.

FIGURE 6 The TCF of $(\text{Bi}_{0.75}\text{Ce}_{0.25})\text{VO}_4$ and $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramics as a function of sintering temperature.

FIGURE 7 Thermal expansion data of the $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic.

Tables

Table 1 Refined atomic fractional coordinates from XRD data of the $(\text{Bi}_{0.05}\text{Ce}_{0.95})\text{VO}_4$ ceramic and the cell parameters are $a = b = 7.3777(4) \text{ \AA}$, $c = 6.4809(1) \text{ \AA}$ with a space group $I4_1/amd$ (No. 141).

Atom	Wyckoff position	Occ.	x	y	z	Biso
Bi	4a	0.05000	0.00000	0.75000	0.12500	0.00914
Ce	4a	0.95000	0.00000	0.75000	0.12500	0.00914
V	4b	1.00000	0.00000	0.25000	0.37500	0.00992
O	16h	1.00000	0.00000	0.07309	0.20880	0.01314

Table 2 Phonon parameters obtained from the fitting of the infrared reflectivity spectra of (Bi_{0.05}Ce_{0.95})VO₄ ceramic.

Mode	ω_{oj}	ω_{pj}	γ_j	$\Delta\epsilon_j$
1	158.60	340.24	29.10	4.60
2	241.10	249.53	39.76	1.07
3	463.28	332.93	161.81	0.52
4	622.37	402.31	120.87	0.42
5	716.04	416.28	67.75	0.34
6	761.22	411.21	37.79	0.29
7	845.58	108.77	47.00	0.02
(Bi _{0.05} Ce _{0.95})VO ₄	$\epsilon_\infty = 2.23$		$\epsilon_0 = 9.48$	











